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# Generation of Picosecond Laser Pulses at 1064 nm from All Solid-State Passively Mode-Locked Lasers

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**Abstract:** Using semiconductor saturable absorber mirror and diode end-pumping configurations, all solid-state passively mode-locked  $Nd^{3+}$ :YVO<sub>4</sub> lasers have been successfully developed. The lasers efficiently provide a stable train of ultra-short laser pulses of 12 ps at 1064 nm at the pulse repetition rate adjustable from 8.8 MHz to 100 MHz. The peak power of 1.3 kW and an average laser power of 940 mW were obtained at the pulse repetitive rate of 60 MHz, corresponding to a laser conversion efficiency of 43%. In order to obtain lower pulse repetitive rate and, therefore, higher peak power, a long laser resonator (larger than 15 m) is proposed and successfully developed for diode end-pumped passively mode-locked laser operation using a simple multiple-pass cavity configuration. As a result, the peak laser power up to 5.1 kW was obtained at the pulse repetitive rate of 8.8 MHz. The experimental results of the picosecond laser amplification and the harmonic generations at 532 nm (2nd), 355 nm (3rd) and 266 nm (4th) are presented.

Key words: laser diode pumping, passive mode-locking, solid-state laser.

# I. INTRODUCTION

Ultrashort pulsed lasers have been widely applied for the time-resolved measurements in physics, chemistry and biology. By the progresses in directly diode-pumped solidstate laser material and by the appearance of a semiconductor saturable absorber mirror (SESAM) mode-lockers, compact ultrafast SESAM mode-locked all-solid laser systems have been strongly developed and dominating the ultrashort solid-state laser market for scientific, technological and medical applications [1-4]. This paper presents the results obtained in the development of CW diode-pumped ultrashort solid-state Nd:YVO<sub>4</sub> lasers passively modelocked with a semiconductor saturable absorber mirror (SESAM). The parameters of SESAM, resonator, optical element and diode laser end-pumping configurations have been researched for stable, highly efficient, passively modelocked laser operations, using commercially available components. The diode-pumped passively mode-locked solid-state Nd:YVO<sub>4</sub> laser generates a stable train of ultrashort laser pulses (12 ps, 8-100 MHz at 1064 nm).

# **II. EXPERIMENTS**

#### **II.1.** Cavity configuration

For a generation of ultrashort laser pulses based on the mode locking techniques, two basic conditions have to be satisfied. Firstly, a laser must generate a fairly large number of longitudinal modes. Secondly, these modes must be equidistant in frequency and synchronized in phase.

In order to satisfy the first condition, a nonselective and sufficiently long laser resonator has been developed for the simultaneous excitation of a sufficiently large number of longitudinal modes. In this case, the multi-frequency generation is very important. Furthermore, the interference of m longitudinal modes which are equidistant in frequency and phase locked leads the laser emission to be the nature of a train of ultrashort pulses which follow one another with time interval T equal to the resonator roundtrip time

We have chosen the laser medium of Nd:YVO<sub>4</sub> that is very efficient and robust as well as easily provided by a number of suppliers. The crystal has a gain line width of about  $10^{12}$  Hz corresponding to minimum obtainable pulse duration of a few picoseconds. The resonator configuration of laser diode-pumped mode-locked Nd:YVO<sub>4</sub> laser is shown in Fig. 1.



Fig. 1. Cavity configuration of the laser diode-pumped passively mode-locked Nd:YVO<sub>4</sub> laser

The pump source is a CW diode laser (ATC, Russia) emitted at the wavelength of around 808 nm with a variable output power from 0 to 2100 mW. Active cooling and temperature stabilization is provided by a built-in Peltier cooling device. The polarization of the laser light is horizontal. The diode has a built-in cylindrical micro lens for its fast axis collimation. This allows us to use simple pump optics, consisting of a single lens of 25 mm focus length. The rear mirror  $M_1$  is a dichroic flat mirror (R > 99% @1064 nm and R < 5% @808 nm). The crystal is a  $3 \times 3 \times 3$  mm, a-cut, 1% doped AR/AR coated Nd:YVO4 (Casix) and mounted on a copper heat sink. The Nd:YVO4 was oriented for maximum absorption at 808 nm and titled by about 4-5 degrees with the optical axis in order to avoid the etalon effect.  $M_2$  is a concave mirror of radius R = 30 cm, it allows to get the best overlap between the pumped volume and the laser mode, and minimize intracavity diffraction losses.  $M_3$ ,  $M_4$  and  $M_5$  are flat mirrors, high reflection at the laser wavelength of 1064 nm.  $M_6$  is a concave mirror with radius r = 15 cm and used to focus the intracavity laser energy to the saturation fluence of the SESAM.

#### Mode-locking laser operation

For phase locking of longitudinal modes, we have used the passive locking method with a semiconductor saturable absorber mirror (SESAM) in the resonator. As presented in Section II.1, the SESAM has a relaxation time much smaller than the resonator round-trip time. Therefore, intracavity laser energy is able to bleach the SESAM at a modulation period of about c/2L (c – light speed in the resonator and L – resonator length) and leads a modelocking laser operation.

The threshold of the CW Nd:YVO<sub>4</sub> laser was obtained at the diode pumping power of 150 mW. When the diode pumping energy is less than 540 mW, the intracavity laser energy was too low compared to the saturation fluence of the SESAM and the laser might run in CW. Increasing the diode pumping energy, the intracavity laser energy increases and reaches to the saturation fluence of the SESAM. This leads the laser behavior to be either Q-switched mode-locking (Fig. 2).



Fig. 2. Q-switch mode-locking operation

Increasing the diode pumping energy more than 540 mW the laser operates in pure mode-locking (Fig. 3). In practice, the pure mode-locking laser operation is quite far above the lasing threshold. A stable mode-locking laser operation was obtained with the diode pumping powers larger than 1500 mW. The pulse energy stability is rms.  $\pm 2\%$ .



Fig. 3. Mode-locking operation

Figures 2 and 3 present the pulse trains obtained from the mode-locked Nd:YVO<sub>4</sub> laser with a fast photodiode (rise time < 200 ps) and a digital oscilloscope (Tektronix 1.5 GHz, 20 GS/s).

# • Power characteristic and pulse width of mode-locking laser

The power characteristic of mode-locking Nd:YVO<sub>4</sub> laser are measured with two different output mirrors (reflectivity of 80% and 94%), as shown in Fig. 4. In mode-locking laser operation, considerably high laser efficiency was achieved. At a maximal pumping power of 2200 mW and the output mirror of 80%, we obtained the laser power of 940 mW corresponding to laser efficiency up to 43%. With the output mirror of 94%, a maximum laser power obtained 660 mW corresponding to laser efficiency of 30%.



Fig. 4. Power characteristic of laser



Fig. 5. Autocorrelation trace of laser pulse

The width of mode-locked laser pulse is measured by an autocorrelation system. The autocorrelation trace is shown in Fig. 5. Assuming that the pulse shape of laser is in form of Gaussian profile so we obtained the pulse width of mode-locking laser is about 12 ps.

# **II.2.** Research and development a diode-pumped passively mode-locked Nd:YVO<sub>4</sub> laser operated at pulse repetition rates lower than 10 MHz

For applications, such as fluorescent lifetime measurement, the fluorescence lifetime of numerous organic molecules is in the range of a nanosecond to hundreds of picoseconds, short laser pulses of several tenths of picoseconds are ideally needed from reliable pump laser sources. The pulse repetition rate is of the pumping laser source used is also an important parameter in experiments: High repetition rates ensure fast acquisition of the fluorescence decay signal and allow the dynamic processes to be studied. If the repetition rate is too high and exceeds several megahertz, problems appear in the complex signal-processing devices during the acquisition of the data. Therefore, a pulse repetition rate less than 10 MHz seems to be a good requirement for this kind of application. However, for modelocked solid state lasers with resonator design in Fig. 1, the repetition rate in the range of 30-100 MHz that determined entirely by their resonator length. To obtain a pulse repetition rate lower than 10 MHz, the resonator length should be larger than 15 m. To achieve such long resonator, one introduced a multiple-pass cavity (MPC) based on the classical Herriott-style MPC with two concave mirrors [5]. This resonator configuration was recently used successfully to decrease the frequency of a Ti:Al<sub>2</sub>O<sub>3</sub> mode-locked laser to 4 MHz [6, 7]. The configuration of the MPC used in our system is in fact a version of the Herriott-style MPC folded in two with the help of two plane mirrors (Fig. 6). The distance of both concave mirrors (r = 2 m) from the two plane mirrors is fixed at approximately the focal length of the concave mirrors. The beam is periodically focused and defocused after each reflection on the concave and plane mirrors, respectively.



Fig. 6. Schematic of the long cavity diode-pumped passively mode-locked Nd:YVO<sub>4</sub> laser

During the alignment of the MPC inside the laser resonator, various lengths of the laser resonator could be used, resulting each time in the operation of the modelocked laser system on a number of different repetition rates. The easiest for optics element arrange and control of the number of reflections inside the MPC is with 6 times reflection in MPC so we are archived with cavity length more than 15 m (corresponding repetition rate about 8.8 MHz). With the use of an output mirror of 20% and 6% transmission at 1064 nm, the laser provides the pulse duration of 12 ps (Fig. 8) and an average laser power of 410 mW, corresponding to about 4 kW peak power (Fig. 7).



Fig. 7. Power characteristic of laser



Fig. 8. Interference auto-correlation trace of pulse duration of 12 ps (assuming a Gaussian profile)

The pulse trains obtained from the mode-locked Nd:YVO<sub>4</sub> laser with a fast photodiode (rise time < 200ps) and a digital oscilloscope (Tektronix 1.5 GHz, 20 GS/s) was shown in Fig. 9.



Fig. 9. presents a pulse train stable operation at repetition rate of 8.8 MHz or 114 ns of pulse separation

# II.3. Laser amplification and wavelength conversion

For the amplifier (Fig. 10), we choose a Nd:YVO<sub>4</sub> crystal (1% doped  $3 \times 3 \times 3$  mm<sup>3</sup>) as a gain medium because of its high emission cross section, in order to have a good energy extraction. To decrease the thermal effects, the Nd:YVO<sub>4</sub> crystal was attached to copper blocks.



Fig. 10 Experimental setup for amplifier and converter wavelength of laser

Optical elements of system includes: (1) CW laser diode – (2) pump lens – (4) Nd:  $YVO_4$  crystal – (3), (5-8), (10-12) and (14) mirrors – (9) KTP crystal – (13) 2-nd half wave plate – (15) lens, (16) 3rd frequency crystal

Average laser powers of 650 mW and 900 mW are then obtained with a 2-pass or 4-pass configuration, respectively (Fig. 11). Increased efficiency of the amplifier results from the pump of laser diode. However, the efficiency of the 4-pass amplifier is not as high as expected; this might be due to the thermal effects.



Fig. 11. Power characteristics of the two and four-pass laser amplifier

For the nonlinear processes, we first use a 5 mm long type I KTP crystal for second harmonic generation. The output green power at 532 nm is up to 130 mW with a pulse duration lightly shortened by non-linear conversion of around 10 ps (Fig. 12). For direct fourth generation, a 5 mm



Fig. 12. Average power at 532 nm in function of input average power at 1064 nm

a 5 mm long BBO crystal was used to the second harmonic generation pulse at 532 nm, the output power at 266 nm could be up to 15 mW. For third harmonic generation, the IR and the green are then recombined in a second non-linear crystal (5 mm long type II KDP). An average power of 45 mW at 8.8 MHz with pulse duration of around 10 ps is obtained at 355 nm, corresponding to a conversion efficiency of 5% from IR to UV. It is noted that the non-linear conversion efficient is straightforward improved much better, due to the use of non-linear crystals and optical components given in our laboratory.

# **III. CONCLUSION**

The research and development of a passively modelocked diode-pumped Nd:YVO<sub>4</sub> laser operating at 1064 nm at pulse repetition rates in range of 30-100 MHz have been achieved. The mode-locking of the laser was obtained with a low-modulation semiconductor saturable absorber mirror (2%). In addition, a super long cavity with multiple-pass resonator configuration folded with two concave mirrors (r = 2 m) is achieved. We also have successfully developed an amplifier laser power and converter wavelength to visible and UV. As a result, the laser system was compact and well arranged on an optical table (L120 × W40 cm<sup>2</sup>). Such a picoseconds laser source is demanded and very attractive for applications in biology and photonics.

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# References

- Principles of lasers. (Ed. O. Svelto. 4th edition, Plenum Press) p. 359 (1998).
- [2] S. Forget, F. Balembois, G. Lucas-Leclin, P. Georges, Opt. Com. 220, 187 (2003).
- [3] N.H. Schiller, X.M. Zhao et al, Appl. Opt. 28, 946 (1989).
- [4] J.B. Deaton, Jr., A.D.W. Mckie, J.B. Spicer, J.W. Wagner, Appl. Phys. Lett. 56, 2390 (1990).
- [5] A. Herriott, H. Kogelnik, R. Kompfner, Appl. Opt. 3, 523 (1994).
- [6] S.H. Cho, F.X. Kärtner, U. Morgner, E.P. Ippen et al., Opt. Lett. 26, 560 (2000).
- [7] S.H. Cho, B.E. Bouma, E.P. Ippen, J.G. Fujimoto, Opt, Lett. 24, 417 (1999).
- [8] U. Keller, D.A. B. Miller, G.D. Boyd, T.H. Chiu et al., Opt. Lett. 17, 505 (1992).
- [9] U. Keller, K.J. Weingarten et al., IEEE J. Sel. Top. Quant. Electron. 2, 453 (1996).
- [10] C. Hönninger, R. Paschotta et al., J. Opt. Soc. Am. B 16, 46 (1999).
- [11] G.J. Spühler, T. Südmeyer, R. Paschotta et al. Appl. Phys. B 71, 19 (2000).
- [12] N.T. Nghia, Do Q. Khanh et al. ASEAN J. Scien.&Tech. for Develop. 24, 1-2, 139 (2007).



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